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B.Sc (Part-II) Hons

Langevin's theory of paramagnetism:-

According to Langevin, each atom has permanent magnetic moment  $m$ . The only force acting on the atom is due to external field  $\vec{B}$ . Let  $\theta$  be the angle of inclination of the axis of the atomic dipole with the direction of the applied field  $\vec{B}$ . Then the magnetic potential energy of the atomic dipole is

$$U = -mB \cos \theta$$

Now, on classical statistics, the number of atoms making an angle between  $\theta$  and  $\theta + d\theta$  is

$$dn = C e^{mB \cos \theta / KT} \sin \theta d\theta$$

Where  $K$  is Boltzmann's Constant and  $T$  is the absolute temperature

putting  $mB/KT = \alpha$

then

~~$$dn = C e^{\alpha \cos \theta} \sin \theta d\theta$$~~

$$dn = C e^{\alpha \cos \theta} \sin \theta d\theta \quad \text{--- (1)}$$

Hence the total number of atomic magnets in unit volume of the

Paramagnetic material

$$n = \int_0^\pi dn$$

$$n = \int_0^\pi c e^{\alpha \cos \theta} \sin \theta d\theta \quad \text{--- (2)}$$

Let us

$$\cos \theta = x$$

$$\therefore dx = -\sin \theta \cdot d\theta$$

$$\text{for } \theta = 0^\circ, \cos \theta = x, x = +1$$

$$\text{for } \theta = \pi, \cos \theta = x, x = -1$$

$$\therefore n = - \int_{+1}^{-1} c e^{\alpha x} dx$$

$$= -c \int_{+1}^{-1} e^{\alpha x} dx$$

$$= -c \left[ \frac{e^{\alpha x}}{\alpha} \right]_{+1}^{-1}$$

$$n = -c/\alpha (e^{-\alpha} - e^{\alpha})$$

$$n\alpha = c(e^{\alpha} - e^{-\alpha})$$

$$\therefore c = \frac{n\alpha}{e^{\alpha} - e^{-\alpha}} \quad \text{--- (3)}$$

The components of each dipole moment parallel to B is  $m \cos \theta$ . The total magnetic moment of all the  $n$  atoms

Contained in unit volume of the gas is the magnetisation  $M$ .

$$\begin{aligned} \therefore M &= \int_0^\pi m \cos \theta \, d\Omega \\ &= \int_0^\pi m \cos \theta \cdot c e^{\chi \cos \theta} \cdot \sin \theta \, d\theta \end{aligned}$$

Again let

$$\chi = \cos \theta$$

$$\therefore d\chi = -\sin \theta \, d\theta$$

$$\therefore M = \int_0^\pi m \cos \theta \cdot c e^{\chi \cos \theta} \cdot \sin \theta \, d\theta$$

$$= - \int_{+1}^{-1} m \chi c e^{\chi \chi} \, d\chi$$

$$= -cm \int_{+1}^{-1} \chi e^{\chi \chi} \, d\chi$$

$$= -cm \left[ \chi \int e^{\chi \chi} \, d\chi - \int \left( \frac{d}{d\chi} (\chi) \int e^{\chi \chi} \, d\chi \right) d\chi \right]_{+1}^{-1}$$

$$= -cm \left[ \chi \frac{e^{\chi \chi}}{\alpha} - \int \frac{e^{\chi \chi}}{\alpha} \, d\chi \right]_{+1}^{-1}$$

$$= -cm \left[ \chi \cdot \frac{e^{\chi \chi}}{\alpha} - \frac{1}{\alpha} \cdot \frac{e^{\chi \chi}}{\alpha} \right]_{+1}^{-1}$$

$$= \frac{-cm}{\alpha} \left[ -e^{-\alpha} - \frac{e^{-\alpha}}{\alpha} - e^{\alpha} + \frac{e^{\alpha}}{\alpha} \right]$$

$$M = \frac{-cm}{\alpha} \left[ \frac{1}{\alpha} (e^{\alpha} - e^{-\alpha}) - (e^{\alpha} + e^{-\alpha}) \right]$$

putting the value of  $c$ , we get

$$M = \frac{mn}{e^\alpha - e^{-\alpha}} \left[ (e^\alpha + e^{-\alpha}) - \frac{1}{\alpha} (e^\alpha - e^{-\alpha}) \right]$$

$$\therefore M = mn \left[ \frac{e^\alpha + e^{-\alpha}}{e^\alpha - e^{-\alpha}} - \frac{1}{\alpha} \right]$$

$$M = mn \left[ \coth \alpha - \frac{1}{\alpha} \right] \quad \text{--- (4)}$$

$$\text{or, } M = mn L(\alpha) \quad \text{--- (5)}$$

where  $L(\alpha) = \left[ \coth \alpha - \frac{1}{\alpha} \right]$  is called the Langevin function

The variation of  $M$  with  $\alpha$  has shown this figure

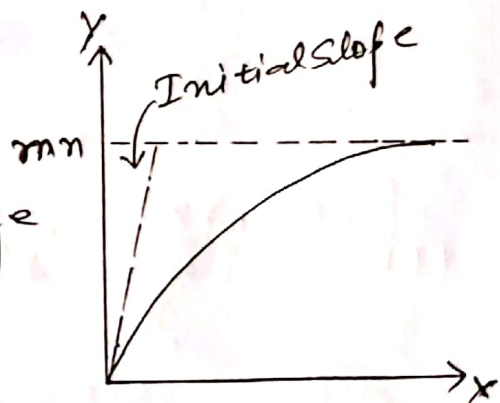
Case-1 :- At low temperature and large applied field,  
 $L(\alpha) \rightarrow 1$

$$\therefore M = mn$$

Thus the saturation point is reached when all the atomic dipoles are parallel to  $\vec{B}$ .

Case-1-11 :- Under normal condition  $\alpha$  is very small,

$$\text{then } L(\alpha) = \coth \alpha - \frac{1}{\alpha} = \frac{\alpha}{3}$$





$$\therefore M = \frac{nm\mu}{3}$$

$$= \frac{nm^2 B}{3kT}$$

$$= \frac{nm^2 \mu_0 H}{3kT} \quad (\because B = \mu_0 H)$$

$$\therefore \chi = \frac{M}{H}$$

$$= \frac{nm^2 \mu_0 H}{3kT \cdot H}$$

$$= \frac{nm^2 \mu_0}{3kT}$$

$$= \frac{\mu_0 nm^2}{3k} \left( \frac{1}{T} \right)$$

$$= C \left( \frac{1}{T} \right) \quad \left( \text{where } C = \frac{\mu_0 nm^2}{3k} = \text{constant} \right)$$

This constant is called Curie constant

$$\therefore \chi = \frac{C}{T} \quad \text{--- (6)}$$

$\chi \propto \frac{1}{T}$   
According to Weiss' modification

$$\chi = \frac{C}{T - \theta}$$

where  $\theta$  is called Curie temperature, which is a characteristic of the substance.